Quantum Mechanics, Physics 531 Homework Assignment 7, due May 9, 2008

Problem 1. Problem 9.1.

Problem 2. Problem 9.15.

Problem 3. Problem 10.9.

Problem 4. Problem 29 in Chapter 5 of Sakurai.

$$\left(-\frac{iE_b}{\hbar}\right)e^{-iE_bt/\hbar}$$
].

I the last two terms

$$r_b e^{-iE_b t/\hbar}$$
. [9.8]

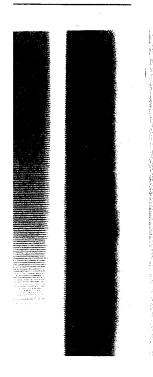
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tion the middle possitum—does not occur).

roceed by spontaneous eme of allowed transigure 9.6. Note that the no lower-energy state is indeed much longer -1). Metastable states -1) alled forbidden

9.74. Hint: First show

trate that



*Problem 9.1 A hydrogen atom is placed in a (time-dependent) electric field $E = E(t)\hat{k}$. Calculate all four matrix elements H'_{ij} of the perturbation H' = eEz between the ground state (n = 1) and the (quadruply degenerate) first excited states (n = 2). Also show that $H'_{ii} = 0$ for all five states. Note: There is only one integral to be done here, if you exploit oddness with respect to z; only one of the n = 2 states is "accessible" from the ground state by a perturbation of this form, and therefore the system functions as a two-state configuration—assuming transitions to higher excited states can be ignored.

**Problem 9.15 Develop time-dependent perturbation theory for a multilevel system, starting with the generalization of Equations 9.1 and 9.2:

$$H_0\psi_n = E_n\psi_n, \quad \langle \psi_n | \psi_m \rangle = \delta_{nm}.$$
 [9.79]

At time t = 0 we turn on a perturbation H'(t), so that the total Hamiltonian is

$$H = H_0 + H'(t). [9.80]$$

(a) Generalize Equation 9.6 to read

$$\Psi(t) = \sum c_n(t)\psi_n e^{-iE_nt/\hbar}, \qquad [9.81]$$

and show that

$$\dot{c}_m = -\frac{i}{\hbar} \sum_n c_n H'_{mn} e^{i(E_m - E_n)t/\hbar},$$
 [9.82]

where

$$H'_{mn} \equiv \langle \psi_m | H' | \psi_n \rangle. \tag{9.83}$$

(b) If the system starts out in the state ψ_N , show that (in first-order perturbation theory)

$$c_N(t) \cong 1 - \frac{i}{\hbar} \int_0^t H'_{NN}(t') dt'.$$
 [9.84]

and

$$c_m(t) \cong -\frac{i}{\hbar} \int_0^t H'_{mN}(t') e^{i(E_m - E_N)t'/\hbar} dt', \quad (m \neq N).$$
 [9.85]

(c) For example, suppose H' is *constant* (except that it was turned on at t = 0, and switched off again at some later time t). Find the probability of transition from state N to state M ($M \neq N$), as a function of t. Answer:

$$4|H'_{MN}|^2 \frac{\sin^2[(E_N - E_M)t/2\hbar]}{(E_N - E_M)^2}.$$
 [9.86]

(d) Now suppose H' is a sinusoidal function of time: $H' = V \cos(\omega t)$. Making the usual assumptions, show that transitions occur only to states with energy $E_M = E_N \pm \hbar \omega$, and the transition probability is

$$P_{N\to M} = |V_{MN}|^2 \frac{\sin^2[(E_N - E_M \pm \hbar\omega)t/2\hbar]}{(E_N - E_M \pm \hbar\omega)^2}.$$
 [9.87]

(e) Suppose a multilevel system is immersed in incoherent electromagnetic radiation. Using Section 9.2.3 as a guide, show that the transition rate for stimulated emission is given by the same formula (Equation 9.47) as for a two-level system.

* *Problem 10.9 Suppose the one-dimensional natificial oscillator (mass in, itequency ω) is subjected to a driving force of the form $F(t) = m\omega^2 f(t)$, where f(t)is some specified function (I have factored out $m\omega^2$ for notational convenience; f(t) has the dimensions of length). The Hamiltonian is

$$H(t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m\omega^2 x^2 - m\omega^2 x f(t).$$
 [10.90]

Assume that the force was first turned on at time t = 0: f(t) = 0 for $t \le 0$. This system can be solved exactly, both in classical mechanics and in quantum mechanics.²¹

(a) Determine the classical position of the oscillator, assuming it started from rest at the origin $(x_c(0) = \dot{x}_c(0) = 0)$. Answer:

$$x_c(t) = \omega \int_0^t f(t') \sin[\omega(t - t')] dt'. \qquad [10.91]$$

(b) Show that the solution to the (time-dependent) Schrödinger equation for this oscillator, assuming it started out in the nth state of the undriven oscillator $(\Psi(x,0) = \psi_n(x))$ where $\psi_n(x)$ is given by Equation 2.61), can be written

$$\Psi(x,t) = \psi_n(x-x_c)e^{\frac{i}{\hbar}\left[-(n+\frac{1}{2})\hbar\omega t + m\dot{x}_c(x-\frac{x_c}{2}) + \frac{m\omega^2}{2}\int_0^t f(t')x_c(t')dt'\right]}.$$
 [10.92]

(c) Show that the eigenfunctions and eigenvalues of H(t) are

$$\psi_n(x,t) = \psi_n(x-f); \quad E_n(t) = \left(n + \frac{1}{2}\right)\hbar\omega - \frac{1}{2}m\omega^2 f^2.$$
 [10.93]

- (d) Show that in the adiabatic approximation the classical position (Equation 10.91) reduces to $x_c(t) \cong f(t)$. State the precise criterion for adiabaticity, in this context, as a constraint on the time derivative of f. Hint: Write $\sin[\omega(t-t')]$ as $(1/\omega)(d/dt')\cos[\omega(t-t')]$ and use integration by parts.
 - (e) Confirm the adiabatic theorem for this example, by using the results in (c) and (d) to show that

$$\Psi(x,t) \cong \psi_n(x,t)e^{i\theta_n(t)}e^{i\gamma_n(t)}.$$
 [10.94]

Check that the dynamic phase has the correct form (Equation 10.39). Is the geometric phase what you would expect?

9. Consider a composite system made up of two spin $\frac{1}{2}$ objects. For t < 0, the Hamiltonian does not depend on spin and can be taken to be zero by suitably adjusting the energy scale. For t > 0, the Hamiltonian is given by

$$H = \left(\frac{4\Delta}{\hbar^2}\right) \mathbf{S}_1 \cdot \mathbf{S}_2.$$

Suppose the system is in $|+-\rangle$ for $t \le 0$. Find, as a function of time, the probability for being found in each of the following states $|++\rangle$, $|+-\rangle$, $|-+\rangle$, and $|--\rangle$:

a. By solving the problem exactly.

b. By solving the problem assuming the validity of first-order timedependent perturbation theory with H as a perturbation switched on at t = 0. Under what condition does (b) give the correct results?